

## ▣ PARTICLE IN AN INFINITE WELL

- Consider the following equation

$$\nabla^2 \psi(x,y) = -\alpha \psi(x,y)$$

where  $\alpha > 0$ . Can we solve this w/ the boundary conditions  $\psi(0,y) = \psi(L,y) = \psi(x,0) = \psi(x,W) = 0$ ?

- Sure, let's start w/ S.o.V. & look for solutions  $f(x)g(y)$ :

$$g(y) \frac{\partial^2 f(x)}{\partial x^2} + f(x) \frac{\partial^2 g(y)}{\partial y^2} = -\alpha f(x)g(y)$$

$$\hookrightarrow \frac{1}{f(x)} \frac{\partial^2 f(x)}{\partial x^2} + \frac{1}{g(y)} \frac{\partial^2 g(y)}{\partial y^2} = -\alpha$$

- Now, I listed two homogeneous boundary conditions for both  $x$  and  $y$ , which suggests trig functions for both variables.
- We couldn't do this for an eqn. like  $\nabla^2 u(x,y) = 0$ , where trig functions in  $x$  imply exponentials in  $y$ , and vice-versa.
- But when I apply S.o.V. to  $\nabla^2 \psi = -\alpha \psi$  this is fine:

$$\frac{1}{f} \frac{\partial^2 f}{\partial x^2} = -k^2 \quad \frac{1}{g} \frac{\partial^2 g}{\partial y^2} = -p^2 \quad -k^2 - p^2 = -\alpha$$

That  $-\alpha$  ( $\alpha > 0$ ) term on the RHS means I can use a negative S.o.V. constant for both my  $x$  &  $y$  eqns.

$$\hookrightarrow f(x) = A \cos(kx) + B \sin(kx) \quad g(y) = C \cos(py) + D \sin(py)$$

$$\text{w/ } k^2 + p^2 = \alpha^2$$

- The B.C. @  $x=0$  tells us  $A=0$ , and the B.C. @  $x=L$  forces  $k = n\pi/L$ ,  $n \in \mathbb{Z}$  &  $n \geq 1$ .
- Likewise, the B.C. @  $y=0$  &  $y=W$  set  $C=0$  &  $p = m\pi/W$  w/  $m \in \mathbb{Z}$  &  $m \geq 1$ .
- So our S.o.V. solutions look like

$$f(x)g(y) = B \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi y}{W}\right) \quad \text{w/} \quad \frac{n^2\pi^2}{L^2} + \frac{m^2\pi^2}{W^2} = \alpha.$$

- Now, depending on the value of  $\alpha$ ,  $L$ , &  $W$  this may or may not work! Suppose I told you  $L=W=1$  and  $\alpha = 25\pi^2$ . Then you'd conclude

$$\frac{n^2\pi^2}{1} + \frac{m^2\pi^2}{1} = 25\pi^2 \Rightarrow n=3 \text{ \& } m=4 \text{ or } n=4 \text{ \& } m=3$$

So in that case there are two solutions.

- But what if I said  $L=1$ ,  $W=2$ , and  $\alpha = 17$

$$\frac{n^2\pi^2}{1} + \frac{m^2\pi^2}{4} = 17 \quad \leftarrow \text{No solutions w/ } n, m \in \mathbb{Z}!$$

In that case our problem has no solutions!

- Maybe a better question is what values of  $\alpha$  let this problem have interesting solutions? In that case the answer is just

$$\alpha = \frac{n^2\pi^2}{L^2} + \frac{m^2\pi^2}{W^2} \quad n, m \in \mathbb{Z} \text{ \& } \geq 1$$

which combines  $k = n\pi/L$  &  $p = m\pi/W$  from our boundary conditions w/ the requirement  $-k^2 - p^2 = -\alpha$  to make S.o.V. work!

- What does this have to do w/ particles and wells?
- In Quantum Mechanics, a particle w/ mass  $M$  that can move in the  $x$  &  $y$  directions w/ potential energy  $U(x,y)$  is described by a WAVE FUNCTION  $\Psi(x,y,t)$  that satisfies the Schrödinger Equation:

$$i\hbar \frac{\partial}{\partial t} \Psi(x,y,t) = \left( -\frac{\hbar^2}{2M} \nabla^2 + U(x,y) \right) \Psi(x,y,t)$$

- What if we used S.o.V. to look for sol'ns of the form  $T(t)\Psi(x,y)$ ? We'd arrive at:

$$\frac{1}{T(t)} i\hbar \frac{\partial T(t)}{\partial t} = \frac{1}{\Psi(x,y)} \left( -\frac{\hbar^2}{2M} \nabla^2 + U(x,y) \right) \Psi(x,y)$$

- Both sides must equal a constant for this to work. Let's call that constant 'E'

$$i\hbar \frac{\partial T}{\partial t} = ET \Rightarrow T(t) = (\text{Const}) e^{-i\frac{E}{\hbar}t}$$

$$-\frac{\hbar^2}{2M} \nabla^2 \Psi(x,y) + U(x,y) \Psi(x,y) = E \Psi(x,y)$$

- When  $E$  is a positive, real const. we call this a 'Stationary state,' since  $\Psi^*(x,y,t)\Psi(x,y,t)$  doesn't change over time. ( $\Psi^*\Psi$  is the thing we do with a W.F. to find the probability that the particle is @ position  $x,y$  @ time  $t$ !)
- This constant  $E$  is just the energy.
- Now imagine a potential energy  $U(x,y)$  that is zero when  $0 \leq x \leq L$  and  $0 \leq y \leq W$ , & infinite everywhere else.

- This is called an 'infinite well' because the particle is trapped in the region  $0 \leq x \leq L$  &  $0 \leq y \leq W$ . No amount of energy lets it 'climb out' since  $U = \infty$  outside the well.
- For the wave function this means  $\Psi(x, y, t) = 0$  if  $x \leq 0$ ,  $x \geq L$ ,  $y \leq 0$ , or  $y \geq W$ , since the particle can never be found there.
- So, inside the  $\infty$ -well, the Schrödinger eqn says a stationary state w/ energy  $E$  satisfies

$$E \Psi(x, y) = -\frac{\hbar^2}{2M} \nabla^2 \Psi(x, y)$$

$$w/\Psi(0, y) = \Psi(L, y) = \Psi(x, 0) = \Psi(x, W) = 0$$

- This is the problem we just solved, with  $\alpha = \frac{2ME}{\hbar^2}$ !
- Since  $E$  was a constant we introduced via S.o.V., our question about "which  $\alpha$  work?" now makes more sense.
- A particle of mass  $M$  in an  $\infty$ -well of length  $L$  & width  $W$  is only allowed to have certain energies. They are:

$$E_{n,m} = \frac{\hbar^2 \pi^2}{2M} \times \left( \frac{n^2}{L^2} + \frac{m^2}{W^2} \right) \quad n, m \in \mathbb{Z} \text{ & } \geq 1$$

The allowed energies are labeled by a pair of integers  $n$  &  $m$ .

The wave function for a state w/ energy  $E_{n,m}$  is:

$$\Psi(x, y, t) = (\text{Constant}) e^{-i \frac{E_{n,m}}{\hbar} t} \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi y}{W}\right)$$

Fixed by requiring  $\int_0^L dx \int_0^W dy |\Psi(x, y, t)|^2 = 1$

- A few notes about this. First,  $\hbar$  has units of J.s, so  $E_{n,m}/\hbar$  indeed has units of (angular) frequency & we often write  $\omega_{n,m} = E_{n,m}/\hbar$  or

$$\omega_{n,m} = \frac{\hbar \pi^2}{2M} \left( \frac{n^2}{L^2} + \frac{m^2}{W^2} \right)$$

- The W.F. times its 'Complex Conjugate' is related to the probability of detecting the particle in the vicinity of  $(x,y)$  around time  $t$

$$P(x,y,t) \propto \Psi(x,y,t) \Psi^*(x,y,t)$$

Here, this means  $\Psi$   
w/  $i \rightarrow -i$

Since the particle has to be somewhere in the well, the probabilities have to add up to '1', and this tells us how to Normalize the W.F.

$$1 = \int_0^L dx \int_0^W dy \Psi(x,y,t) \Psi^*(x,y,t)$$

So if the system is in a state w/ energy  $E_{n,m}$

$$\begin{aligned} 1 &= (\text{const})^2 \int_0^L dx \int_0^W dy e^{-i \frac{E_{n,m}}{\hbar} t} \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi y}{W}\right) e^{+i \frac{E_{n,m}}{\hbar} t} \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi y}{W}\right) \\ &= (\text{const})^2 \underbrace{\int_0^L dx \sin^2\left(\frac{n\pi x}{L}\right)}_{L/2} \underbrace{\int_0^W dy \sin^2\left(\frac{m\pi y}{W}\right)}_{W/2} \end{aligned}$$

$$\hookrightarrow (\text{const})^2 = \frac{4}{LW} \Rightarrow (\text{const}) = \frac{2}{\sqrt{LW}}$$

So the normalized W.F. for the stationary state  $(n,m)$  is

$$\Psi(x,y,t) = \frac{2}{\sqrt{LW}} e^{-i \frac{E_{n,m}}{\hbar} t} \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi y}{W}\right)$$

$$\omega/E_{n,m} = \frac{\hbar^2 \pi^2}{2M} \times \left( \frac{n^2}{L^2} + \frac{m^2}{W^2} \right)$$